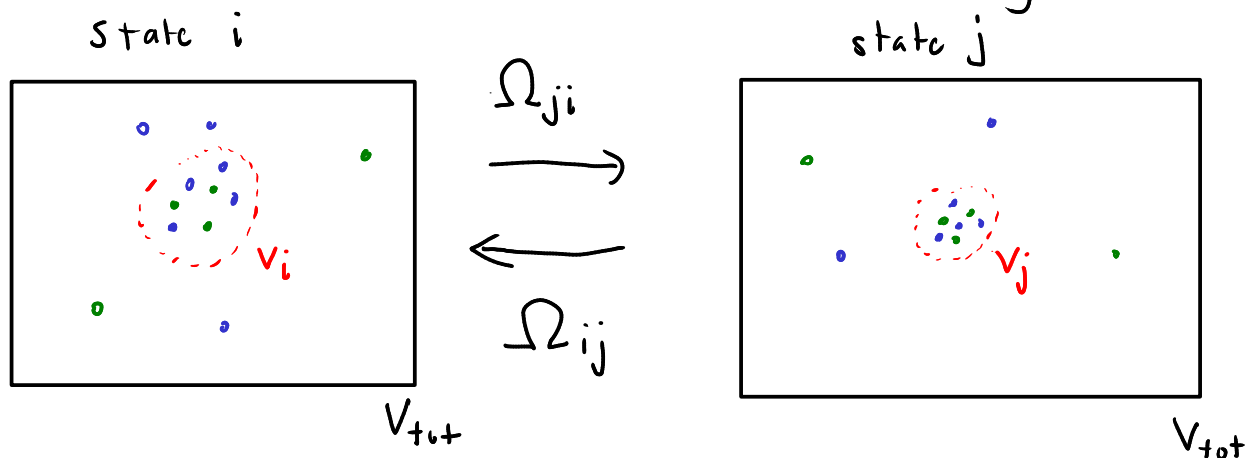


model where certain types of molecules cannot leave / enter cell (on regular timescales) but cell volume can change (b/c of water moving in / out)



$n_\alpha = \# \text{ molec. of type } \alpha \text{ inside cell}$
 $N_\alpha - n_\alpha = \# \text{ molec. of type } \alpha \text{ outside cell}$

focus on impermeable molec. here

n_α doesn't change

$$\frac{\Omega_{ji}}{\Omega_{ij}} = e^{-\beta(-k_B T \ln \frac{\delta_j}{\delta_i})}$$

$\delta_j = \# \text{ microstates in state } j$

$\delta_i = \# \text{ microstates in state } i$

$N_{pos, \alpha}^{in, i} = \# \text{ pos. available to molec. of type } \alpha \text{ on inside of cell in state } i$

$$= \frac{V_i}{V_{mol, \alpha}}$$

\uparrow vol. of type α molec.

$$N_{pos, \alpha}^{out, i} = \frac{V_{tot} - V_i}{V_{mol, \alpha}}$$

$$\sigma_i \approx \underbrace{\begin{pmatrix} N_{\text{pos}, \alpha}^{\text{in}, i} \\ n_\alpha \end{pmatrix} \begin{pmatrix} N_{\text{pos}, \alpha}^{\text{out}, i} \\ N_\alpha - n_\alpha \end{pmatrix}}_{\text{type } \alpha \text{ configs}} \underbrace{\begin{pmatrix} \phantom{N_{\text{pos}, \alpha}^{\text{in}, i}} \\ \phantom{N_{\text{pos}, \alpha}^{\text{out}, i}} \end{pmatrix} \begin{pmatrix} \phantom{N_{\text{pos}, \alpha}^{\text{in}, i}} \\ \phantom{N_{\text{pos}, \alpha}^{\text{out}, i}} \end{pmatrix}}_{\text{similar factors for all other types}} \circ \circ \circ$$

"dilute approx." = ignore volume of other types when calc. config's of type α

$$\ln \binom{M}{m} \approx -m \ln \frac{m}{M} \quad m \gg 1, M \gg 1$$

$$M \gg m$$

$$\Rightarrow \ln \sigma_i \approx - \sum_{\substack{\alpha \\ \in \text{types}}} \left[n_\alpha \ln \frac{n_\alpha V_{\text{mol}, \alpha}}{V_i} + (N_\alpha - n_\alpha) \cdot \ln \frac{(N_\alpha - n_\alpha) V_{\text{mol}, \alpha}}{V_{\text{tot}} - V_i} \right]$$

recall: $\frac{\Omega_{ji}}{\Omega_{ij}} = e^{-\beta(E_j - E_i)} \xrightarrow{t \rightarrow \infty} p_i^s = \frac{e^{-\beta E_i}}{Z}$

here: $\frac{\Omega_{ji}}{\Omega_{ij}} = e^{-\beta(-k_B T \ln \sigma_j - (-k_B T \ln \sigma_i))}$

$$\xrightarrow{t \rightarrow \infty} p_i^s = \frac{e^{-\beta(-k_B T \ln \sigma_i)}}{Z} = \frac{e^{\ln \sigma_i}}{Z}$$

Z ensures $\sum_i p_i^s = 1$

most likely state i (highest p_i^s)

\Rightarrow need largest $\ln \delta_i$

$$\frac{d}{dV_i} \ln \delta_i = 0 \quad \xrightarrow{\text{algebra}} \quad \sum_{\alpha} \left[\underbrace{-\frac{n_{\alpha}}{V_i}}_{C_{\alpha}^{\text{in}}} + \underbrace{\frac{(N_{\alpha} - n_{\alpha})}{V_{\text{tot}} - V_i}}_{C_{\alpha}^{\text{out}}} \right] = 0$$

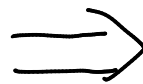
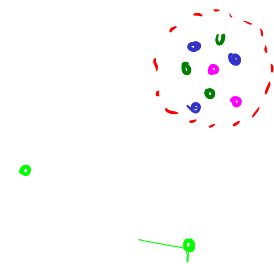
most likely state characterized by:

$$\Rightarrow \sum_{\alpha} C_{\alpha}^{\text{in}} = \sum_{\alpha} C_{\alpha}^{\text{out}}$$

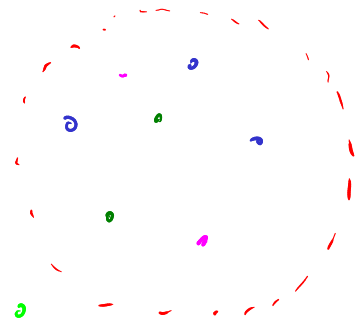
inside
conc. of
 α molec.

outside
conc. of
 α molec.

sum of all inside conc. of impermeable molec. = outside conc. of impermeable molec.



water flows in as system goes to equil.



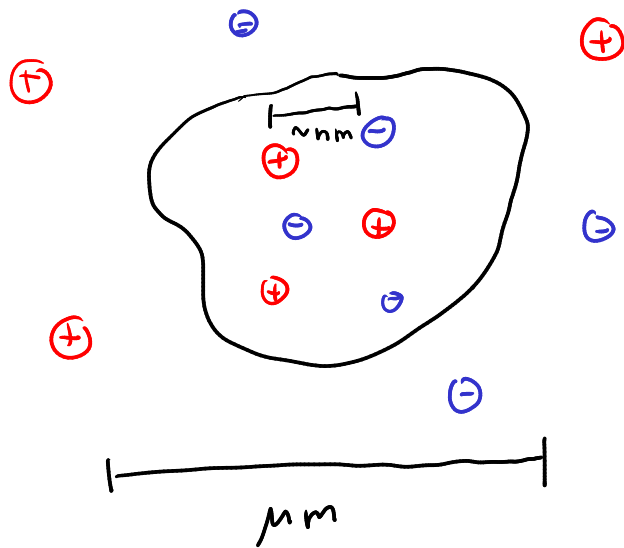
osmotic stress

this can lead to cell exploding!

all cells w/ flexible membranes must have transporters or other mechanisms to regulate osmotic stress

III. Charged particles + transport across membranes

Starting point:



in a cell
many molecules are
charged but
there is an approx.
balance of pos.
& neg. charges
on inside & outside
 \approx electroneutrality

typically: few nm b/t charges

pair of $+e$ & $-e$ ions: Coulomb energy of attraction

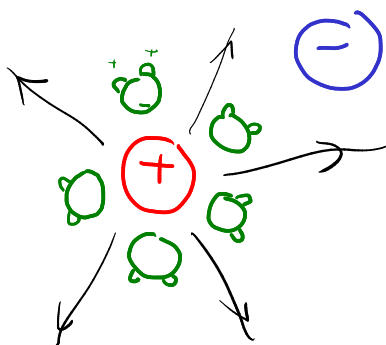
$r = \text{separation} \sim 1 \text{ nm}$

$$= \frac{e^2}{4\pi\epsilon r}$$

$$\epsilon = \epsilon_r \epsilon_0$$

\downarrow perm. of surrounding material
 \downarrow ϵ_0 vacuum permittivity
 \downarrow ϵ_r relative permittivity

 $\left\{ \begin{array}{l} \approx 1 \text{ air} \\ \approx 80 \text{ water} \end{array} \right.$



Water reorientation
weakens overall field \Rightarrow large ϵ_r
seen by $-e$ charge

$r = 1 \text{ nm}$, $\epsilon_r = 80 \Rightarrow$ energy $\sim -0.7 k_B T$

\Rightarrow weak attraction easily broken by thermal fluctuations of order $\sim k_B T$

\Rightarrow keeps ions dispersed in solution