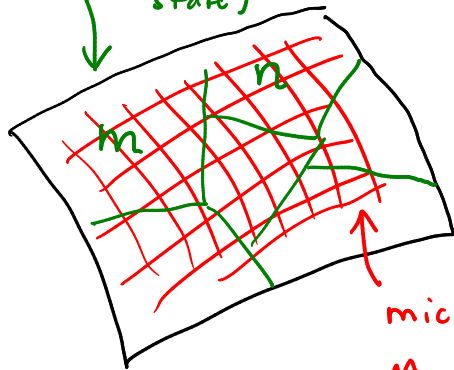


Summary: 1) isolated ^{total} system w/ const total energy E_{tot} which is ergodic + mixing

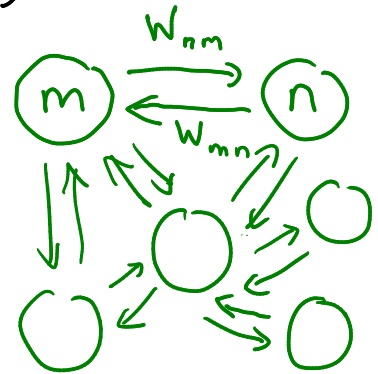
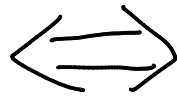
2) divide up total into "system" (focus) + "environment" (rest)

macrostates
(describe system state)



const.
 E_{tot}

microstates
 m (const. phase volume = a)



3) network of transitions w/ matrix W whose entries satisfy local detailed balance (LDB):

$$P_n^s = \frac{\sum_{\mu} \Omega_{\mu} \rightarrow n}{\sum_{\mu} \Omega_{\mu}} \quad \# \text{ microst. in macrost. } n$$

$$W_{nm} P_m^s = W_{mn} P_n^s$$

Consequences of LDB:

1) microscopic reversibility (MR):

$$\text{if } W_{nm} \neq 0 \Rightarrow W_{mn} \neq 0$$

if arrow exists in one direc., has to exist in reverse

\Rightarrow get a unique stationary state

2) Can we describe this stationary state?

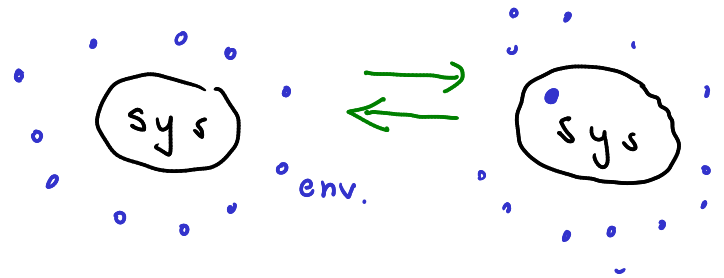
To do this, we need a couple more assumptions:

- Ω_n = counts # of microstates available to environ. when sys. is in state n & hence should depend on conserved quantities describing environ.

examples: energy $E_{\text{tot}} = E_n + E_n^{\text{env}} \Rightarrow E_n^{\text{env}} = E_{\text{tot}} - E_n$

sys env.

exchange of particles



$$N_{\text{tot}}^{\text{part}} = N_n^{\text{part}} + N_n^{\text{part,env}}$$

part. in tot. sys env.

Ω_n is a function of conserved quantities:

$$\Omega_n = \Omega(E_n^{\text{env}}, N_n^{\text{part,env}}, \dots)$$

↑ universal function (\approx continuous)

focus on case where only energy is exchanged:

LDB: $W_{nm} p_m^s = W_{mn} p_n^s$

\uparrow $\frac{\Omega_m}{\Omega_{\text{tot}}}$ \uparrow $\frac{\Omega_n}{\Omega_{\text{tot}}}$

$$\Rightarrow \frac{W_{nm}}{W_{mn}} = \frac{\Omega_n}{\Omega_m} = \frac{\Omega(E_{\text{tot}} - E_n)}{\Omega(E_{\text{tot}} - E_m)}$$

env \gg system: $E_{tot} \gg E_n$ Taylor expand!

$$\frac{W_{nm}}{W_{mn}} = \exp \left[\ln \Theta(E_{tot} - E_n) - \ln \Theta(E_{tot} - E_m) \right]$$

$$\ln \Theta(E_{tot} - E_n) \approx \ln \Theta(E_{tot}) - \underbrace{\frac{\partial \ln \Theta}{\partial E}}_{\beta} \Big|_{E_{tot}} E_n + \dots$$

$$= \ln \Theta(E_{tot}) - \beta E_n + \dots$$

$$\Rightarrow \frac{W_{nm}}{W_{mn}} = \exp \left[-\beta (E_n - E_m) + \dots \right]$$

traditionally: $\beta \equiv \frac{1}{k_B T} = \frac{\partial \ln \Theta}{\partial E} \Big|_{E_{tot}}$

definition of temp. T $k_B =$ Boltzmann's const.
 $T =$ units of Kelvins (K) $= 1.38 \times 10^{-23} \text{ J/K}$

$k_B T =$ units of J

$\beta =$ units of J^{-1}

uphill $= \frac{W_{nm}}{W_{mn}} = e^{-\beta(E_n - E_m)}$ \swarrow how "easy" it is to get energy from env. case:
 downhill $\quad \quad \quad \underbrace{\hspace{2cm}}_{\text{energy diff.}} \quad \begin{array}{c} E_n \\ \uparrow W_{nm} \quad \downarrow W_{mn} \\ E_m \end{array} \quad \begin{array}{l} E_n > E_m \\ \beta > 0 \end{array}$
 $=$ needs to be donated from env. if uphill trans.

$\beta > 0$ ($T > 0$): uphill \ll downhill } environ. is "stingy": less likely to give than to take away energy
 (typical for env \gg sys)

$\beta \rightarrow +\infty$ ($T \rightarrow 0^+$): $\frac{\text{uphill}}{\text{downhill}} \rightarrow 0$ all uphill trans. are forbidden
(absolute zero temp.)

$\beta \rightarrow 0^+$ ($T \rightarrow +\infty$): $\frac{\text{uphill}}{\text{downhill}} \rightarrow 1$ uphill trans. equally likely to downhill

untypical case:

$\beta < 0$ ($T < 0$): $\frac{\text{uphill}}{\text{downhill}} > 1$ extremely generous environ.

\Rightarrow how do we get this?

next time: example of a "gas" of spins

$\ln \Theta(E_{\text{tot}})$

