

$$\Omega_n = \Omega(E_n^{\text{env}}, N_n^{\text{part, env}}, \dots)$$

# microstates  
in macrostate n

focus on case where only energy is conserved.

$$\text{LDB} = W_{nm} p_m^S = W_{mn} p_n^S$$

$$\uparrow \qquad \qquad \qquad \uparrow$$

$$\Omega_m / \Omega_{\text{tot}} \qquad \Omega_n / \Omega_{\text{tot}}$$

$$\frac{W_{nm}}{W_{mn}} = \frac{\Omega_n}{\Omega_m} = \frac{\Omega(E_n^{\text{env}})}{\Omega(E_m^{\text{env}})} = \frac{\Omega(E_{\text{tot}} - E_n)}{\Omega(E_{\text{tot}} - E_m)}$$

assume: env  $\gg$  system  $E_{\text{tot}} \gg E_n$

$$\frac{W_{nm}}{W_{mn}} = \exp \left[ \ln \Omega(E_{\text{tot}} - E_n) - \ln \Omega(E_{\text{tot}} - E_m) \right]$$

$$\ln \Omega(E_{\text{tot}} - E_n) \approx \ln \Omega(E_{\text{tot}}) - \underbrace{\frac{\partial \ln \Omega}{\partial E}}_{\beta} \Big|_{E_{\text{tot}}} E_n + \dots$$

$$\approx \ln \Omega(E_{\text{tot}}) - \beta E_n + \dots$$

$$\Rightarrow \frac{W_{nm}}{W_{mn}} = \exp \left[ -\beta (E_n - E_m) + \dots \right]$$

traditionally:  $\beta = \frac{1}{k_B T} = \left. \frac{\partial \ln \Theta}{\partial E} \right|_{E_{\text{tot}}}$

$\beta =$  inverse energy =  $J^{-1}$   
units

$k_B =$  Boltzmann's constant

$T =$  units of Kelvin (K)

$= 1.38 \times 10^{-23} \text{ J/K}$

$k_B T =$  units of J

$\frac{W_{nm}}{W_{mn}} = \frac{\text{uphill}}{\text{downhill}} = e^{-\beta(E_n - E_m)}$

"how easy it is to absorb energy from env."

$E_n > E_m$   
 $\beta > 0$

energy diff. that needs to be absorbed from env.

$\beta > 0$  ( $T > 0$ ):  $\frac{\text{uphill}}{\text{downhill}} < 1$

env. is "stingy":  
less likely to give energy to sys. than take it away

our everyday experience: typical when env  $\gg$  sys

$\beta \rightarrow +\infty$  ( $T \rightarrow 0^+$ )  $\frac{\text{uphill}}{\text{downhill}} \rightarrow 0$

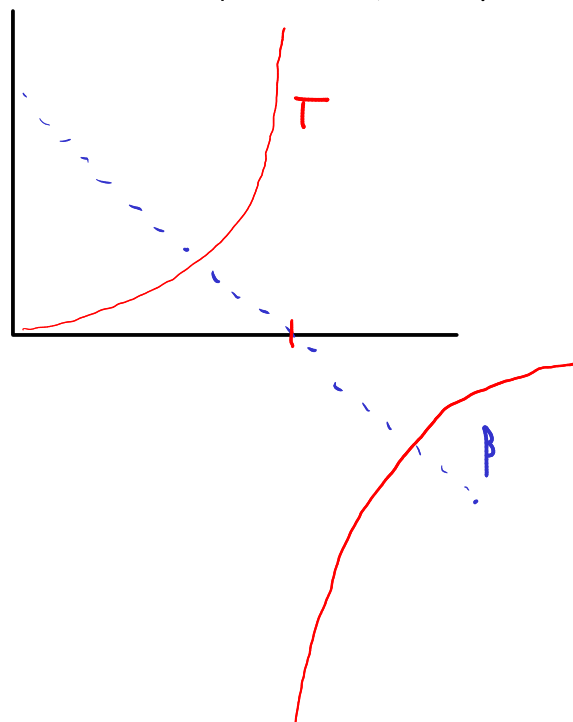
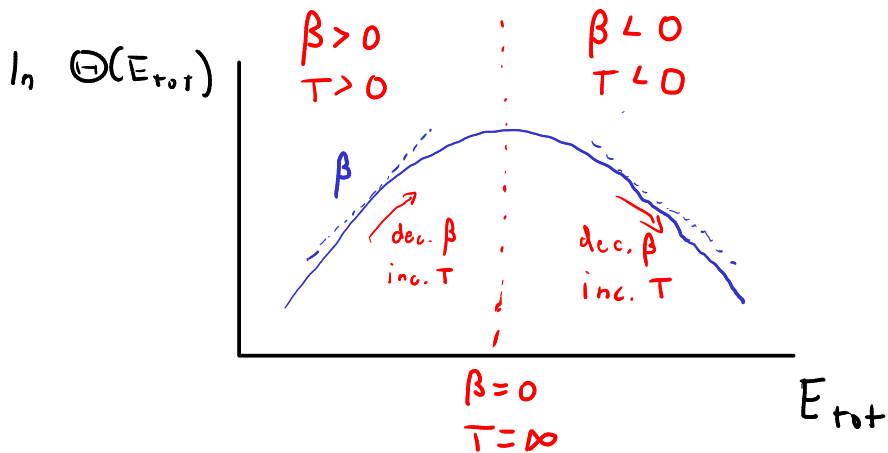
all uphill trans. are forbidden (abs. zero temp.)  
equally likely to gain or lose energy to env.

$\beta \rightarrow 0^+$  ( $T \rightarrow +\infty$ )  $\frac{\text{uphill}}{\text{downhill}} \rightarrow 1$

untypical case: really hot! hotter than  $T = +\infty$

$\beta < 0$  ( $T < 0$ ):  $\frac{\text{uphill}}{\text{downhill}} > 1$

extremely "generous" environment!



Consequences:

$$P_n^s = \frac{\Theta_n}{\Theta_{tot}} = \frac{\Theta_n}{\sum_m \Theta_m}$$

$$\ln \Theta_n = \ln \Theta(E_{tot} - E_n) \approx \ln \Theta(E_{tot}) - \beta E_n$$

$$= \frac{e^{\ln \Theta_n}}{\sum_m \Theta_m}$$

$$= \frac{\Theta(E_{tot})}{\sum_m \Theta_m} e^{-\beta E_n} \Rightarrow P_n^s = \frac{e^{-\beta E_n}}{Z}$$

const. indep. of  $n \equiv \frac{1}{Z}$

$$P_n^s = \frac{e^{-\beta E_n}}{Z}$$

Boltzmann equit. distrib.

note:

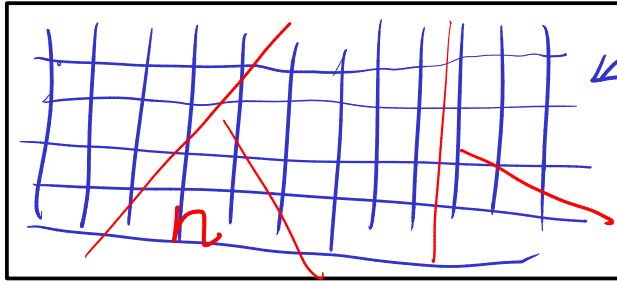
$$\Theta_{tot} \neq \Theta(E_{tot})$$

$$\Theta_{tot} = \sum_m \Theta(E_{tot} - E_m)$$

$Z$  = "partition func" = normaliz. const.

$$\sum_n p_n^s = 1 \Rightarrow Z = \sum_m e^{-\beta E_m} \quad \text{ensemble: microcanonical}$$

ergodic  
+  
mixing



$$P_\mu^s = \frac{1}{\Omega_{\text{tot}}}$$

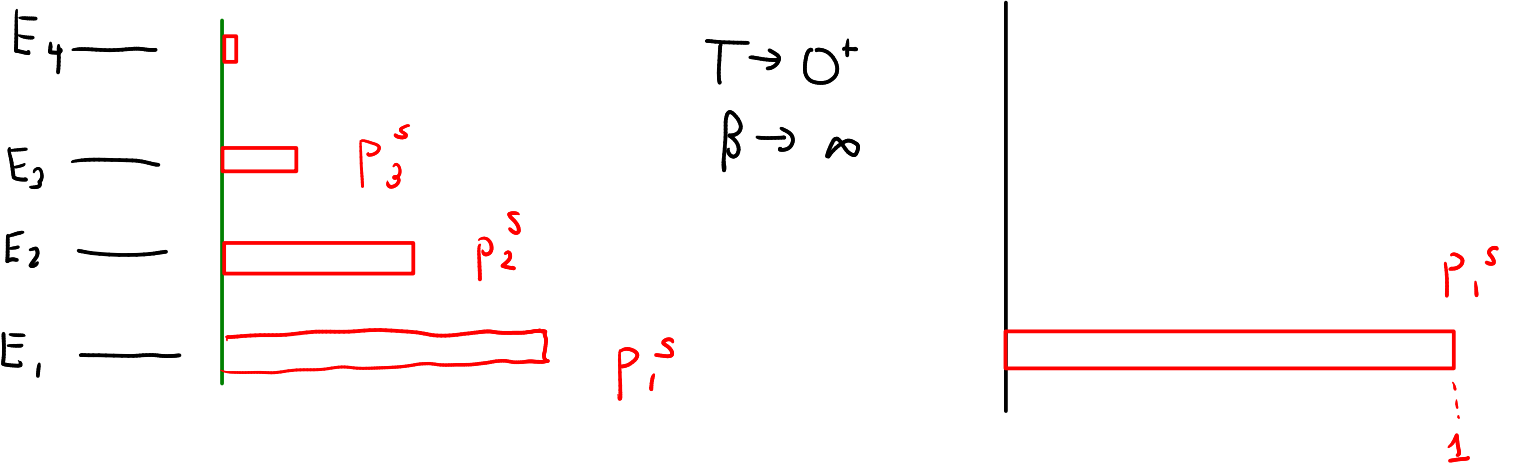
canonical ensemble

$$P_n^s = \frac{e^{-\beta E_n}}{Z}$$

$$T > 0$$

$$\beta > 0$$

⇒ same story, different perspectives



$E_1, E_2, \dots$

$$P_1^s = \frac{e^{-\beta E_1}}{e^{-\beta E_1} + e^{-\beta E_2} + \dots} \xrightarrow{\beta \rightarrow \infty} 1$$

$$P_2^s = \frac{e^{-\beta E_2}}{e^{-\beta E_1} + e^{-\beta E_2} + \dots} \xrightarrow{\beta \rightarrow \infty} e^{-\beta(E_2 - E_1)} \rightarrow 0$$