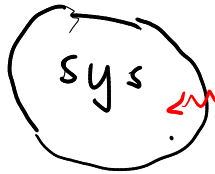


# Modifying LDB condition:

up to now:

env



energy exchange

$$\frac{W_{nm}}{W_{mn}} = e^{-\beta(E_n - E_m)}$$

$$W_{mn} = e^{-\beta Q_{nm}}$$

$Q_{nm}$  = energy taken from therm. "heat" environment during  $m \rightarrow n$  trans.

$> 0$  gain from env.

$< 0$  lost to env.

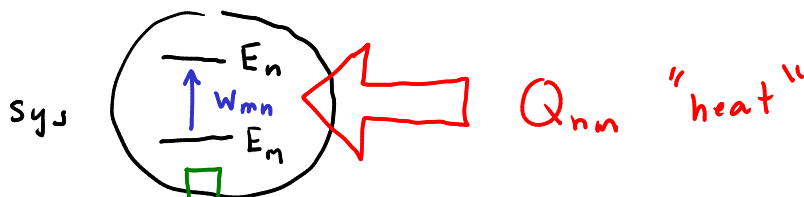
$$\beta = \frac{1}{k_B T}$$

characterizes "generosity" of env.

add a work term for  $m \rightarrow n$  transition:  $W_{nm}$

$W_{nm} > 0$ :

sys. does work on something ext.



$W_{nm}$  = work done by sys. during  $m \rightarrow n$  trans.

$W_{nm} < 0$ : something does work on sys.

$$Q_{nm} = E_n - E_m + W_{nm}$$

need more heat if  $W_{nm} > 0$

generalized version of LDB

$$\frac{W_{nm}}{W_{mn}} = e^{-\beta Q_{nm}}$$

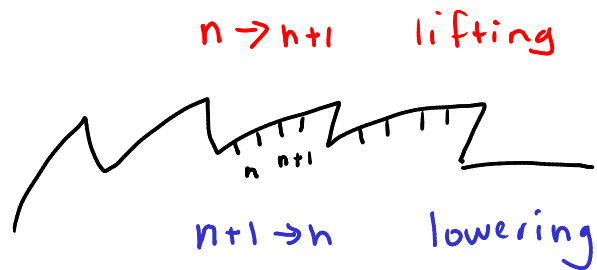
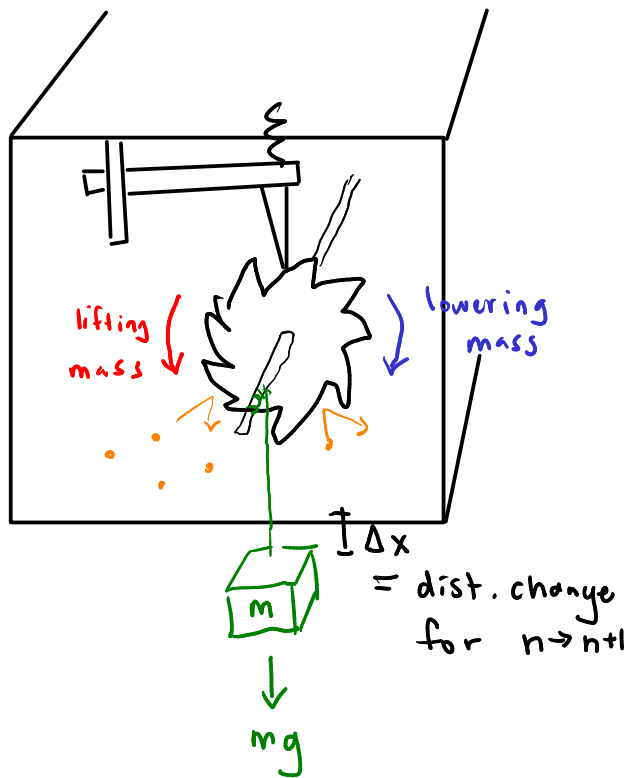
universal form

(just change  $Q_{nm}$   
depending on physical  
model)

$$\frac{W_{nm}}{W_{mn}} = e^{-\beta(E_n - E_m + W_{nm})}$$

↑ depends on trans.  
could be zero

if  $W_{nm} > 0 \Rightarrow$  makes it harder to  
go uphill in energy



$$W_{n+1,n} = mg \Delta x$$

$$W_{n,n+1} = -mg \Delta x$$

**Question:** can we set up a stationary  
state where mass is on avg.  
lifted up (i.e. perpetual motion  
machine)  $\Rightarrow$  return to this later

Return to formalism:  $v = (n_0, n_1, \dots, n_\tau)$

$$I(v) = \sum_{i=1}^{\tau-1} I(\mu_i) \quad \mu_i = (n_i, n_{i+1})$$

$$I(\mu_i) = k_B \ln \frac{W_{n_{i+1}n_i} p_{n_i}(t_i)}{W_{n_i n_{i+1}} p_{n_{i+1}}(t_{i+1})}$$

$$= -k_B \ln p_{n_{i+1}}(t_{i+1}) - (-k_B \ln p_{n_i}(t_i))$$

$$- \frac{1}{T} (E_{n_{i+1}} - E_{n_i}) - \frac{1}{T} W_{n_{i+1}n_i}$$

$$I(v) = \underbrace{-k_B \ln p_{n_\tau}(t_\tau) - (-k_B \ln p_{n_0}(t_0))}_{\equiv \Delta S(v)}$$

"change of entropy"  
in traj.  $v$

$$- \frac{1}{T} \underbrace{(E_{n_\tau} - E_{n_0})}_{\Delta E(v)}$$

$\Delta E(v)$

"change in energy"  
in traj.  $v$

$$- \frac{1}{T} \underbrace{\sum_{i=0}^{\tau-1} W_{n_{i+1}n_i}}_{\bar{W}(v)}$$

$\bar{W}(v)$

"total work" done during  
traj.  $v$

$$\Rightarrow I(v) = \Delta S(v) - \frac{1}{T} \Delta E(v) - \frac{1}{T} \bar{W}(v) \quad \text{defined for every } v$$

notation:  $\Delta A(v) = A_{n_T} - A_{n_0}$

diff. in quantity  $A$   
from initial to final state

$$\langle I(v) \rangle = I$$

$$\langle \Delta A(v) \rangle = \Delta A, \text{ etc.}$$

avg. over ensemble of traj:

$$I = \Delta S - \frac{1}{T} \Delta E - \frac{1}{T} \bar{W} \geq 0$$

we showed  
 $I \geq 0$   
earlier

one version of 2nd  
law of thermodynamics

unpack 
$$\Delta S(v) = -k_B \ln p_{n_T}(t_T) - (-k_B \ln p_{n_0}(t_0))$$

depends on knowing prob's  
in entire ensemble

$$\text{"surprisal"} = -k_B \ln p_n(t_i)$$

large  $p_n(t_i) \Rightarrow$  not surprised to observe  $n$  at time  $t_i$

$\Rightarrow$  low surprisal value  
( $\rightarrow 0$  as  $p_n(t_i) \rightarrow 1$ )

small  $p_n(t_i) \Rightarrow$  very surprised to observe  $n$

$\Rightarrow$  high surprisal

$$\begin{aligned}\Delta S(v) &= \text{surprisal of } n_\tau \text{ at time } t_\tau \\ &\quad - \text{surprisal of } n_0 \text{ at time } t_0 \\ &= \text{"entropy diff. of a traj."}\end{aligned}$$