

case 1: $\dot{Q}_2 > 0 \Rightarrow \eta = \frac{P}{\dot{Q}_2} = 1 - \frac{T_1}{T_2} - T_1 \delta$
 $\Rightarrow \dot{Q}_1 < 0 \leq 1 - \frac{T_1}{T_2} \equiv \eta_{\text{max}}$
 Carnot bound

case 2: $\dot{Q}_1 > 0$ (draw heat out of cold env.)

1st: $P = \dot{Q}_1 + \dot{Q}_2$
 2nd: $\delta = -\frac{\dot{Q}_1}{T_1} - \frac{\dot{Q}_2}{T_2} \geq 0$
 (Stat. state or cycle)

$\Rightarrow \frac{\dot{Q}_2}{T_2} = -\frac{\dot{Q}_1}{T_1} - \delta$
 $\Rightarrow \dot{Q}_2 < 0$
 dumping heat into hot envir.

example: refrigerators
 heat pumps

$\dot{Q}_2 = -\frac{T_2}{T_1} \dot{Q}_1 - T_2 \delta \quad \frac{T_2}{T_1} > 1$
 $\Rightarrow |\dot{Q}_2| > |\dot{Q}_1|$
 $\Rightarrow P = \dot{Q}_1 + \dot{Q}_2 = \dot{Q}_1 - |\dot{Q}_2| < 0$
 \Rightarrow needs net power

input into system
 \Rightarrow plug in your fridge

coeff. of performance $\eta_R = \frac{\dot{Q}_1}{-P} = \frac{\text{heat rate out of cold bath}}{\text{input power}}$

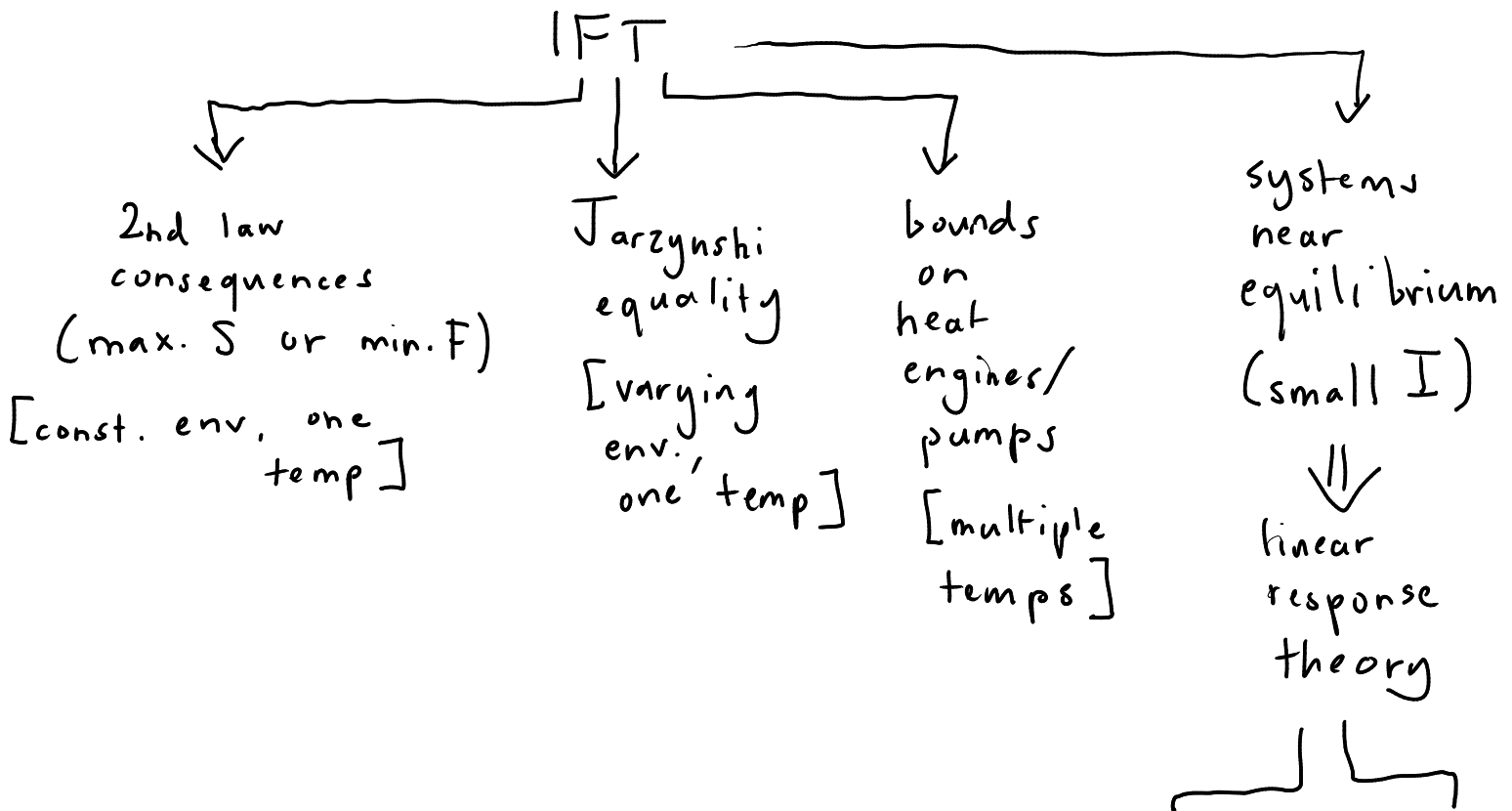
$$\eta_R = \frac{\dot{Q}_1}{-\dot{Q}_1 - \dot{Q}_2} = \text{algebra} = \frac{T_1}{T_2 - T_1} - \underbrace{\frac{T_1 T_2 \delta}{T_2 - T_1}}_{\leq 0}$$

$$\eta_R \leq \frac{T_1}{T_2 - T_1} \equiv \eta_R^{\max}$$

$\sim 300\text{K}$ $\sim 30\text{K}$

η_R^{\max} can be quite large
 (on order ≈ 10)

Big picture perspective: "many children of IFT"



Linear thermodynamics

\downarrow Onsager reciprocity
 \downarrow fluctuation-dissipation theorem

Single temp. T & stat. state
 (not necessarily equil.)

2nd: $I = -\frac{1}{T} \Delta F - \frac{1}{T} W \geq 0$ $\Delta F = \Delta E - T \Delta S$

1st: $Q = \Delta E + W$

$$I = -\frac{1}{T} W = -\frac{1}{T} Q \Rightarrow \begin{matrix} Q = W = -TI \leq 0 \\ \leq 0 & \leq 0 \\ \text{heat} & \text{work} \\ \text{into env.} & \text{on sys.} \end{matrix}$$

$I = 0 \Rightarrow$ ESS ($Q = W = 0$)

$I > 0 \Rightarrow$ NESS ($Q = W < 0$)

"near" equil: I small ($Q = W$ is small)

focus on one time step δt : $\mu_0 = (n_0, n_1)$

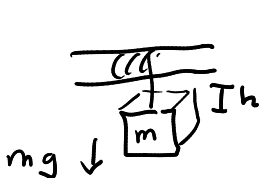
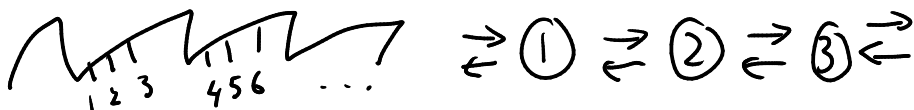
$$I = \langle I(\mu_0) \rangle = -\frac{1}{T} \overline{W} = -\frac{1}{T} \langle W(\mu_0) \rangle$$

\downarrow
 $P_{n_0}^s$

\downarrow
 $P_{n_1}^s$

assume form for work: $\overline{W}(\mu_0) = \omega_{n_1, n_0}$

$= \langle f \Delta x_{n_1, n_0} \rangle$
 \uparrow "force" \uparrow "distance":
 (some phys. quantity we control) change in phys. obser.



$$f = mg$$

$$\Delta x_{n_1, n_0} = \begin{cases} +h & \text{if } n_1 = n_0 + 1 \\ -h & \text{if } n_1 = n_0 - 1 \\ 0 & \text{otherwise} \end{cases}$$

$$f = 0 \Rightarrow \overline{W}(\mu_0) = 0 \Rightarrow \overline{W} = 0 \Rightarrow \overline{I} = 0$$

ESS

assume: $\delta x_{nm} = -\delta x_{mn}$ $\delta x_{nn} = 0$

when $f \neq 0$: $\overline{W} = -f \langle \delta x \rangle = -f \sum_{n, n_0} P(\mu_0) \delta x_{n, n_0}$

$$= -f \sum_{n, n_0} W_{n, n_0} P_{n_0}^s \delta x_{n, n_0}$$

LDB: $\frac{W_{n, n_0}}{W_{n_0, n_1}} = e^{-\beta(E_{n_1} - E_{n_0} - f \delta x_{n, n_0})}$

$$W \vec{p}^s = \vec{p}^s$$

\Rightarrow both W + \vec{p}^s (it's right e-vec)
depend on f

in limit $f \rightarrow 0$: $P_n^s \rightarrow P_n^{eq} = \frac{e^{-\beta E_n}}{Z}$
(ESS)

$W_{nm} \rightarrow W_{nm}^{eq}$ (equil. trans. matrix)

$$\underbrace{W_{nm} P_m^s}_{g(f)} = W_{nm}^{eq} P_m^{eq} \left(1 + c_{nm} f + \dots \right)$$

\uparrow coeff. related to Taylor exp: $\frac{g'(0)}{g(0)}$

$$g(f) = g(0) + g'(0) f + \frac{1}{2} g''(0) f^2 + \dots$$

$$= g(0) \left(1 + \frac{g'(0)}{g(0)} f + \frac{1}{2} \frac{g''(0)}{g(0)} f^2 + \dots \right)$$

plug into avg. \overline{W} expression

$$\begin{aligned}
 W = & -f \sum_{n, n_0} \underbrace{W_{n, n_0}^{eq} p_{n_0}^{eq}}_{\text{Symm. under } n_0 \leftrightarrow n_1} \underbrace{\delta x_{n, n_0}}_{\text{anti-symm.}} \left. \vphantom{\sum_{n, n_0}} \right\} \text{first term} \\
 & -f \sum_{n, n_0} W_{n, n_0}^{eq} p_{n_0}^{eq} C_{n, n_0} f \delta x_{n, n_0} + \dots
 \end{aligned}$$

$= 0$

LDB in equil: $W_{n, n_0}^{eq} p_{n_0}^{eq} = W_{n_0, n_1}^{eq} p_{n_1}^{eq}$