

$$H(\vec{\phi}, \vec{I}) = \sum_{k=1}^N \omega_k I_k$$

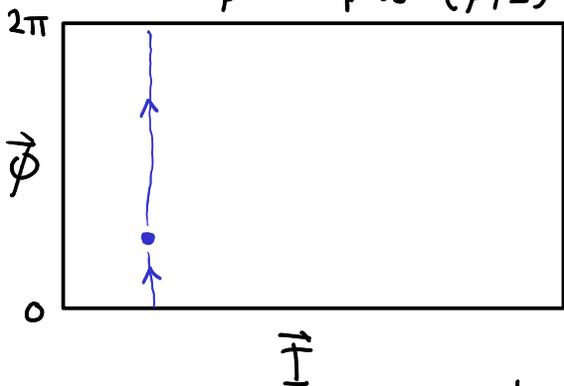
$\underbrace{\omega_k I_k}_{\substack{\text{kth} \\ \text{normal} \\ \text{mode energy}}}$

Hamilton's
equ's
in new coord
phase space $(\vec{\phi}, \vec{I})$

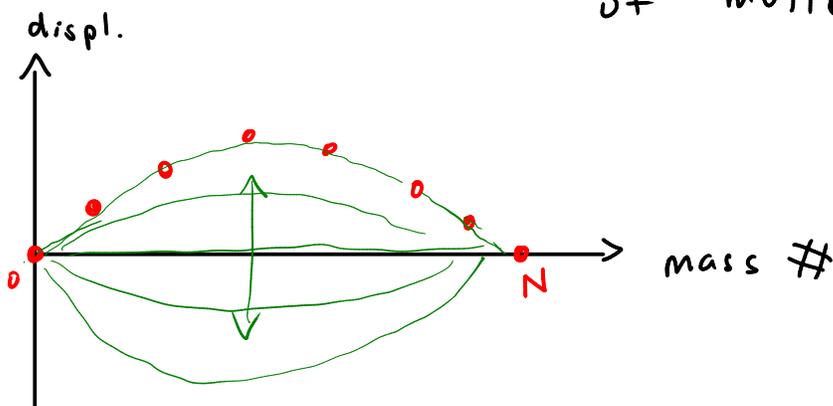
$$\dot{\phi}_i = \frac{\partial H}{\partial I_i} = \omega_i$$

$$\Rightarrow \phi_i(t) = \omega_i t + \phi_i(0)$$

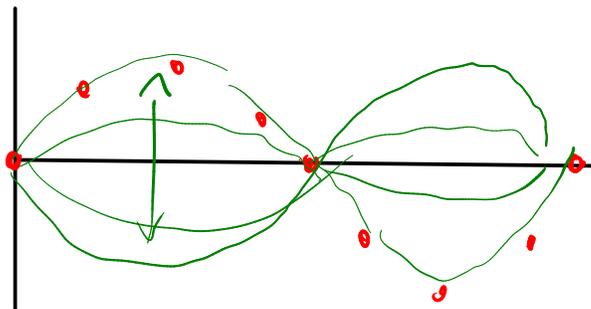
$$\dot{I}_i = -\frac{\partial H}{\partial \phi_i} = 0 \Rightarrow \text{all } I_i \text{ are constants of motion}$$



$k=1$
mode



$k=2$
mode



What happens when $\alpha \neq 0$?

I_N same coords:

$$H(\vec{\phi}, \vec{I}) = \sum_{k=1}^N \omega_k I_k + \alpha U(\vec{\phi}, \vec{I}) + \dots$$

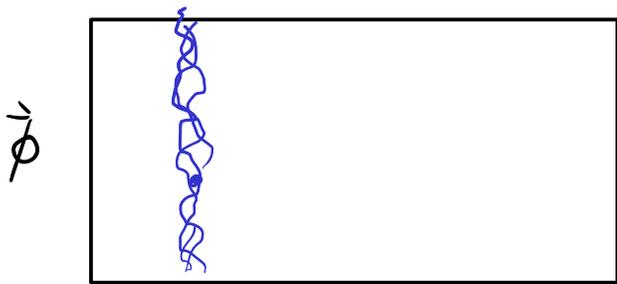
\uparrow
 Small α
 perturbation

\dots
 higher
 order
 in α

idea: $\dot{I}_i = -\frac{\partial H}{\partial \phi_i} \neq 0 \Rightarrow I_i$ are no longer constants of motion

\Rightarrow all of phase space should be explored?

Summary: $\alpha \neq 0$ small FPUT observed quasiperiodic behavior: you return close to (but not exactly) your initial conditions



never achieve microcanonical ensemble (mixing) expected by Fermi

\uparrow stay in narrow band of \vec{I}

Consider a more general class of

systems

$$H(\vec{x}, \vec{p})$$

$$\vec{x} = (x_1, \dots, x_n)$$

$$\vec{p} = (p_1, \dots, p_n)$$

2n-dim
phase
space

This system is integrable:

i) there are n linearly indep. constants of motion: $F_k(\vec{x}, \vec{p}) \quad k=1, \dots, n$

$$\Rightarrow \frac{dF_k}{dt} = \{F_k, H\} = 0 \quad \text{b/c const. of motions}$$

by convention: $F_1 = H$

ii) $\{F_k, F_l\} = 0$ for all k, l

consequence of integrability:

\Rightarrow Liouville-Arnold theorem: if system is integrable, there exists a canonical transf. to "action-angle" coordinates:

$(\vec{\phi}, \vec{I})$
↑ "angles" ↑ "actions"
 $(\phi_1, \dots, \phi_n) \quad (I_1, \dots, I_n)$

note: $\phi_i + 2\pi m \equiv \phi_i$ \downarrow integer

$$H = H(I)$$

$$\dot{I}_k = -\frac{\partial H}{\partial \phi_k} = 0$$

all I_k are const. of motion

$$\phi_k(t) = \omega_k(I)t + \phi_k(0) \quad \Leftarrow \quad \dot{\phi}_k = \frac{\partial H}{\partial I_k}(\vec{I}) \equiv \omega_k(\vec{I})$$

\uparrow const.

examples of integrable systems:

- all 1D problems w/ conserved energy
- n coupled harmonic springs
- central force problems
- two-body grav. problems
- gyroscopes + tops
- free particles confined on surfaces of ellipsoids

not integrable:

- three body problem
(proven by Poincaré to be not integrable)
- chaotic systems
- dissipative systems