

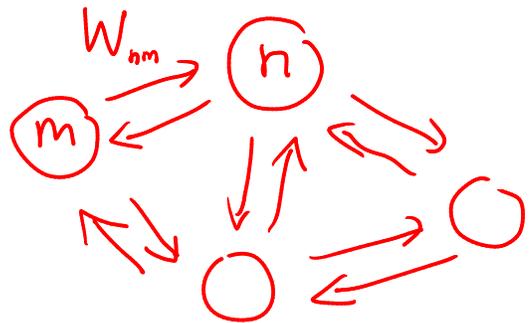
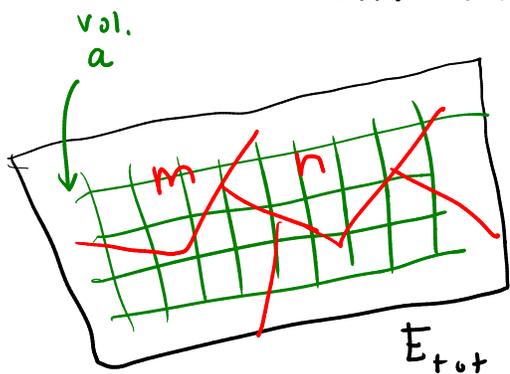
$$\text{LDB: } W_{nm} p_m^s = W_{mn} p_n^s$$

relates elements of W (relev. to transitions at all times) to elements of \vec{p}^s ($t \rightarrow \infty$ limit)

Zoom out:

1) isolated total (system + environment) w/ tot. energy E_{tot} . Assume total is ergodic + mixing.

2) divide up total into system + environment



$$p_n^s = \frac{\textcircled{H}_n}{\textcircled{H}_{tot}} \leftarrow \begin{array}{l} \text{boxes} \\ \text{in } n \end{array}$$

\textcircled{H}_{tot} \leftarrow boxes on whole surface

$$p_m^s = \frac{\textcircled{H}_m}{\textcircled{H}_{tot}}$$

3) network of transitions w/ matrix W satisfies:

$$W_{nm} p_m^s = W_{mn} p_n^s \quad \text{LDB} \quad \equiv$$

consequences of LDB + fact that components of \vec{p}^s have to be nonzero for erg. + mixing systems:

if $W_{nm} \neq 0 \Rightarrow W_{mn} \neq 0$

if arrow in one direction \Rightarrow must be arrow in other

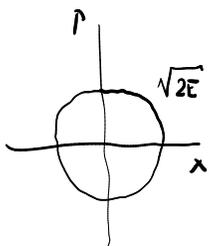
\Rightarrow microscopic reversibility (MR)

\Rightarrow get a unique stationary state

described by $p_n^s = \frac{\Omega_n}{\Omega_{\text{tot}}}$

Properties of this stationary state:

Ω_n = counts # of microstates (boxes) available to environment when sys. is in state n

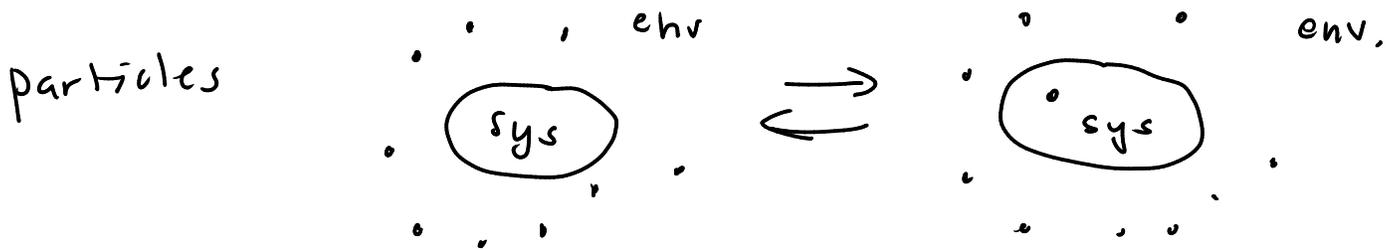


\Rightarrow should depend on energy available to enviro, when sys is in state n

(+ any other conserved quantity b/t sys. + env.)

examples: energy $E_{tot} = E_n^{sys} + E_n^{env}$

$$E_n^{env} = E_{tot} - E_n^{sys}$$



$$N_{tot}^{part} = N_n^{part, sys} + N_n^{part, env}$$

$$\Theta_n = \Theta \left(E_n^{env}, N_n^{part, env}, \dots \right)$$

Some function (universal) other conserved quantities

focus: only energy is exchanged b/t sys. + environ.

$$\Theta_n = \Theta(E_n^{env})$$

$$W_{nm} \frac{\Theta_m^s}{\Theta_{tot}} = W_{mn} \frac{\Theta_n^s}{\Theta_{tot}} \Rightarrow \frac{W_{nm}}{W_{mn}} = \frac{\Theta_n}{\Theta_m} = \frac{\Theta(E_n^{env})}{\Theta(E_m^{env})}$$

$$\Rightarrow \frac{W_{nm}}{W_{mn}} = \frac{\Omega(E_{tot} - E_n)}{\Omega(E_{tot} - E_m)}$$

assume:
env \gg system
 $E_{tot} \gg E_n$

$$\frac{W_{nm}}{W_{mn}} = \exp \left[\ln \Omega(E_{tot} - E_n) - \ln \Omega(E_{tot} - E_m) \right]$$

$$\ln \Omega(E_{tot} - E_n) \approx \ln \Omega(E_{tot}) - \underbrace{\frac{\partial \ln \Omega}{\partial E}}_{\beta} \Big|_{E_{tot}} E_n + \dots$$

$$\approx \ln \Omega(E_{tot}) - \beta E_n$$

$$\frac{W_{nm}}{W_{mn}} = \exp \left[-\beta (E_n - E_m) + \dots \right]$$

LDB in terms of energy
for systems that only
exchange energy w/ env.

traditionally
(def'n of temp.)

$$\beta = \frac{1}{k_B T} = \frac{\partial \ln \Omega}{\partial E} \Big|_{E_{tot}}$$

$$k_B = \text{Boltzmann's const.} \\ = 1.38 \times 10^{-23} \text{ J/K}$$

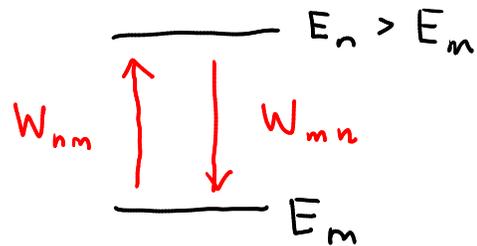
$$\beta = \text{inverse energy} = \text{J}^{-1}$$

$$T = \text{units of Kelvin}$$

$$k_B T = \text{units of J}$$

$$\frac{W_{nm}}{W_{mn}} = \frac{\text{uphill trans.}}{\text{downhill trans.}} = e^{-\beta(E_n - E_m)}$$

energy
diff.
needed
from
env. to
go $m \rightarrow n$
for sys.



$$\beta > 0$$

$\beta > 0$: uphill $<$ downhill
 $W_{nm} < W_{mn}$

higher $\beta \Rightarrow$ harder to go uphill
(lower temp. T)