

$$\frac{W_{nm}}{W_{mn}} = e^{-\beta \underbrace{(E_n - E_m + \omega_{nm})}_{Q_{nm}}} \quad \beta = \frac{1}{k_B T}$$

return to def'n of $I(v)$

$$v = (n_0, \dots, n_\tau)$$

$$\mu_i = (n_i, n_{i+1})$$

$$I(v) = k_B \ln \frac{\mathcal{P}(v)}{\tilde{\mathcal{P}}(\tilde{v})} = \sum_{i=1}^{\tau-1} I(\mu_i)$$

$$I(\mu_i) = k_B \ln \frac{W_{n_{i+1} n_i} p_{n_i}(t_i)}{W_{n_i n_{i+1}} p_{n_{i+1}}(t_{i+1})}$$

$$= -k_B \ln p_{n_{i+1}}(t_{i+1}) - (-k_B \ln p_{n_i}(t_i)) \\ - \frac{1}{T} (E_{n_{i+1}} - E_{n_i}) - \frac{1}{T} \omega_{n_{i+1} n_i}$$

$$\Rightarrow I(v) = -k_B \ln p_{n_\tau}(t_\tau) - (-k_B \ln p_{n_0}(t_0))$$

$$\equiv \Delta S(v)$$

"change in entropy" in traj. v

$$- \frac{1}{T} \underbrace{(E_{n_\tau} - E_{n_0})}_{\Delta E(v)}$$

$$\Delta E(v)$$

"change in energy" in traj. v

$$- \frac{1}{T} \sum_{i=0}^{\tau-1} \omega_{n_{i+1} n_i}$$

$$\} W(v)$$

"total work" done in traj. v

$$\bar{I}(v) = \Delta S(v) - \frac{1}{T} \Delta E(v) - \frac{1}{T} W(v)$$

notation: $\Delta A(v) = A_{n_\tau} - A_{n_0}$

diff in quantity A
from initial to final state in v

$$Q(v)$$

$$\langle Q(v) \rangle = Q$$

$$\langle \Delta A(v) \rangle = \Delta A$$

avg. over ensemble of traj:

$$\boxed{I = \Delta S - \frac{1}{T} \Delta E - \frac{1}{T} W \geq 0}$$

from earlier
conseq.
of
IFT

our version of 2nd
law of thermodyn.

unpack $\Delta S(v) = -k_B \ln p_{n_\tau}(t_\tau)$
 $- (-k_B \ln p_{n_0}(t_0))$

↳ quantity associated w/ particular
traj. v , but it depends on
knowing prob. dist. $\vec{p}(t)$
in whole ensemble

$$\text{"Surprisal"} = -k_B \ln p_n(t_i)$$

if small $p_n(t_i) \Rightarrow$ very surprised to observe state n at time t_i

\Rightarrow high surprisal

if large $p_n(t_i) \Rightarrow$ not surprised

\Rightarrow low surprisal

$$\begin{aligned} \Delta S(v) &= \text{surprisal of } n_\tau \text{ at time } t_\tau \\ &\quad - \text{surprisal of } n_0 \text{ at time } t_0 \\ &= \text{"entropy diff. of a traj."} \end{aligned}$$

unpack work term:

$$W(v) = \sum_{i=0}^{\tau-1} \omega_{n_{i+1}n_i} \stackrel{?}{=} A_{n_\tau} - A_{n_0} \stackrel{\text{can we write?}}{=} \Delta A(v)$$

if this is true: $W(v)$ corresponds to conservative work

traditional def'n in classical mech:

$$\omega = \nabla U \quad \text{there exists a potential } U$$

$$\Rightarrow \underbrace{\int_{\text{path}} \omega}_{\text{total work done over path}} = U(\text{end}) - U(\text{beginning})$$

compare: $W(v) = A_{n_\tau} - A_{n_0}$

if no such $A(v)$ exists $\Rightarrow W(v)$ is not conservative

lastly: heat from environ. during $n_i \rightarrow n_{i+1}$ trans.

energy cons.

$$Q_{n_{i+1}|n_i} = E_{n_{i+1}} - E_{n_i} + \omega_{n_{i+1}|n_i}$$

total heat over traj. $Q(v) = \sum_{i=0}^{\tau-1} Q_{n_{i+1}|n_i} = \Delta E(v) + W(v)$

avg. over ensemble

$$Q = \Delta E + W$$

1st law of thermod.